

Intro to Quantum Computing

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UCL

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Aims

- I can't make you an expert in a day (sorry!)
- Be able to build and analyse quantum circuits
- Understand how quantum computers work
- Understand how quantum computing fits in
- Feel confident to learn more on your own

Outline

- 1 Why bother?
- 2 Computational Logic
- 3 Quantum States
- 4 Quantum Circuits
- 5 Information & Communication
- 6 Hardware
- 7 Compilers

Menti

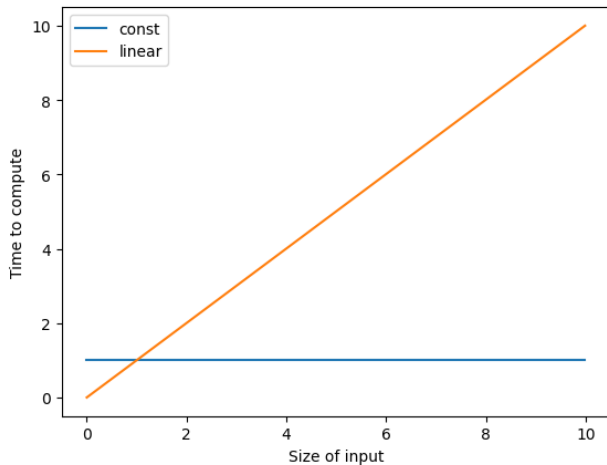
Before we start, join the mentimeter interactive session.



Why bother?

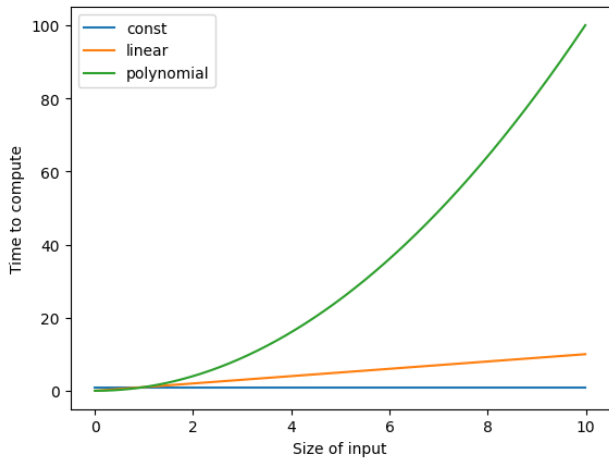
Complexity

- Check odd or even
- Add two numbers



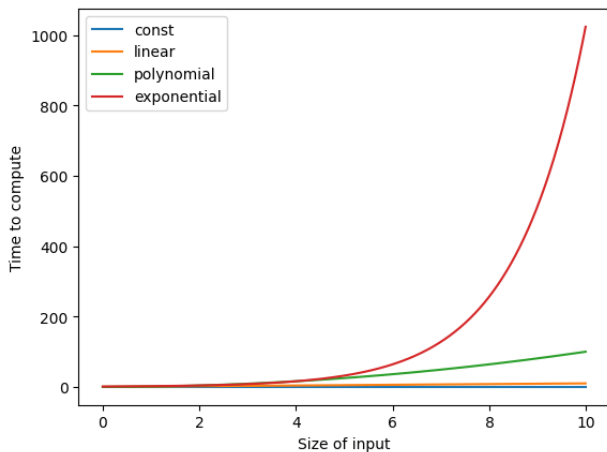
Complexity

- Check odd or even
- Add two numbers
- Multiply two numbers



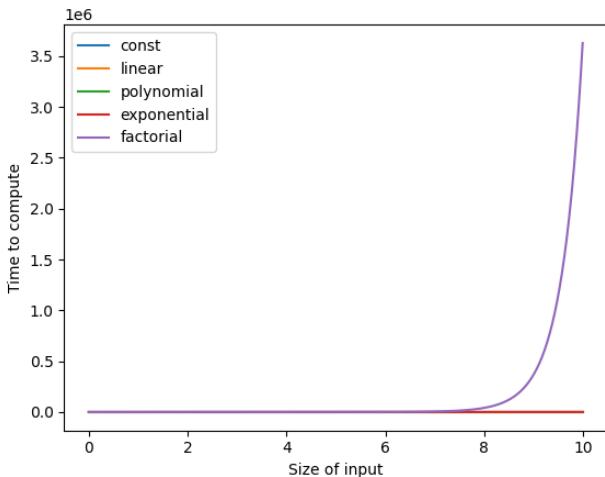
Complexity

- Check odd or even
- Add two numbers
- Multiply two numbers
- Factorise a number



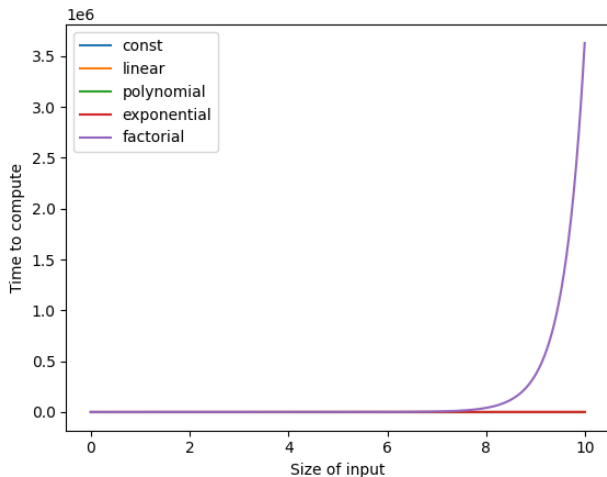
Complexity

- Check odd or even
- Add two numbers
- Multiply two numbers
- Factorise a number
- Find the quickest route for a delivery truck



Complexity

- Find the quickest route for a delivery truck
- Find all possible energy states of a molecule



Feynman

Simulating Physics with Computers

Richard P. Feynman

Department of Physics, California Institute of Technology, Pasadena, California 91107

Received May 7, 1981

1. INTRODUCTION

On the program it says this is a keynote speech—and I don't know what a keynote speech is. I do not intend in any way to suggest what should be in this meeting as a keynote of the subjects or anything like that. I have my own things to say and to talk about and there's no implication that anybody needs to talk about the same thing or anything like it. So what I want to talk about is what Mike Dertouzos suggested that nobody would talk about. I want to talk about the problem of simulating physics with computers and I mean that in a specific way which I am going to explain. The reason for doing this is something that I learned about from Ed Fredkin, and my entire interest in the subject has been inspired by him. It has to do with learning something about the possibilities of computers, and also something about possibilities in physics. If we suppose that we know all

Uses for Quantum Computers

- Design Pharmaceuticals

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- Design Pharmaceuticals
- Save energy

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- Machine Learning

Uses for Quantum Computers

- Design Pharmaceuticals
- Save energy
- Encryption / Decryption
- Machine Learning
- More we don't know yet!

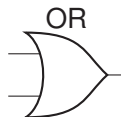
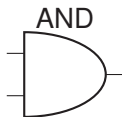
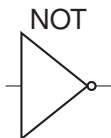
Computational Logic

What does a computer do?

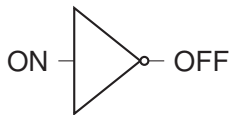
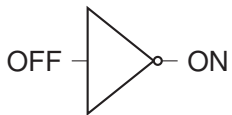
We're going to approach quantum computing from the computing side, and add in quantum mechanics as we go.

- 1 Arithmetic / Logic
- 2 Read / Write Memory
- 3 Control Flow

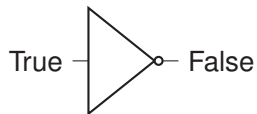
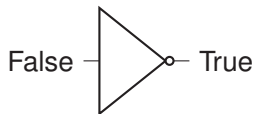
Logic Gates



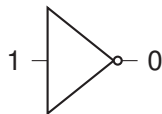
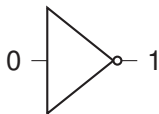
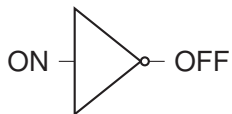
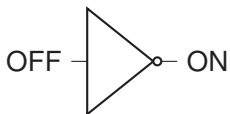
NOT Gate



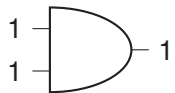
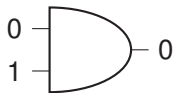
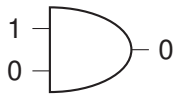
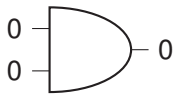
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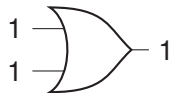
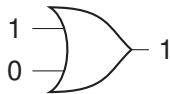
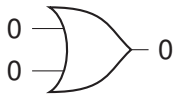
NOT Gate



AND Gate

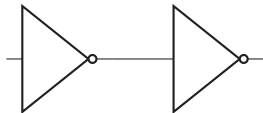


OR Gate



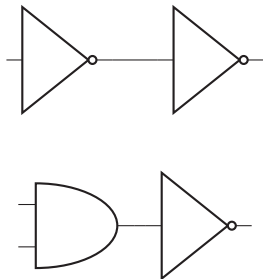
Combining gates

We can combine gates into circuits.



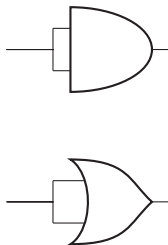
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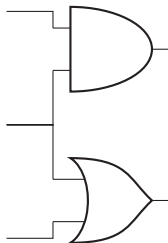
FANOUT

We can split the wire in two.

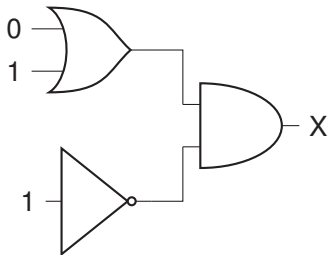


FANOUT

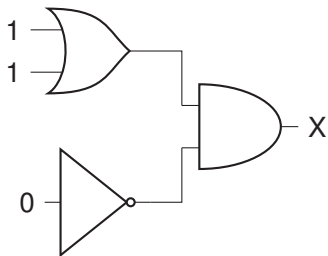
This allows us to share inputs between gates.



First Circuit



First Circuit



Universality of NAND



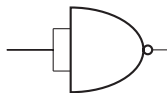
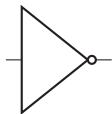
Theorem

All logical operations can be completed with the NAND gate, they are said to be universal.

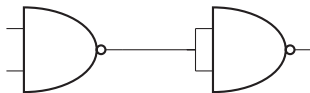
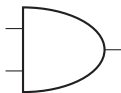
Try to put these together to make something which behaves the same way as NOT, AND and OR gate.¹

¹(Hint: You can split lines or cross them over!)

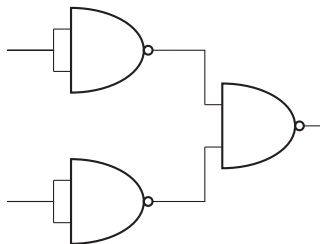
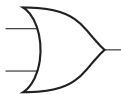
Universality of NAND



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Universality of NAND



States

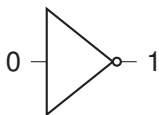
For binary systems we can define two states of each input.

$$\text{OFF} = |0\rangle$$

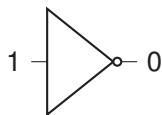
$$\text{ON} = |1\rangle$$

Operators

Rather than draw the gate we can use its name, like we would for a function.



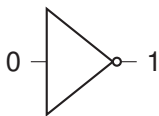
$$\text{NOT } |0\rangle = |1\rangle$$



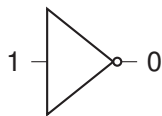
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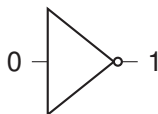
We can go one step further and use a symbol.

$$\hat{N}|0\rangle = |1\rangle$$

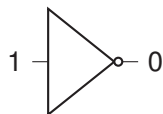
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Gates are represented by symbols with hats, called *operators*.

Multi-bit Operators

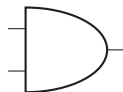
We can apply multi-bit operators to multi-bit states.

$$\hat{A}|00\rangle = |0\rangle$$

$$\hat{A}|01\rangle = |0\rangle$$

$$\hat{A}|10\rangle = |0\rangle$$

$$\hat{A}|11\rangle = |1\rangle$$



Multi-bit Operators

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Rules for Operators

Operators act to their right.

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$$\hat{N}|0\rangle = |1\rangle$$

$$\hat{N}\hat{A}|11\rangle = \hat{N}\left(\hat{A}|11\rangle\right) = \hat{N}|1\rangle = |0\rangle$$

Rules for Operators

Operators only act on states with the correct size of input

$$\hat{A}|0\rangle = 0$$

$$\hat{N}|11\rangle = 0$$

Note that this is 0 and not a state at all!

Rules for Operators

Operators can be applied multiple times.

$$\hat{N}\hat{N}|1\rangle = \hat{N}|0\rangle = |1\rangle$$

Identity Operator

We define the Identity Operator \hat{I} which does nothing to change the state.

$$\hat{I} |0\rangle = |0\rangle$$

$$\hat{I} |1\rangle = |1\rangle$$

This is equivalent to a wire in our circuit.



Identity Operator

This allows us write about operators independent of states.

$$\hat{N}\hat{N}|0\rangle = \hat{N}|1\rangle = \hat{I}|0\rangle$$

$$\hat{N}\hat{N}|1\rangle = \hat{N}|0\rangle = \hat{I}|1\rangle$$

$$\hat{N}\hat{N} = \hat{I}$$

States

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This is essentially the question 'are these the same state?'

Multi-bit States

To combine single bit states into larger sizes we use a *tensor product*.²

$$|x\rangle \otimes |y\rangle = |xy\rangle$$

e.g. $|0\rangle \otimes |1\rangle = |01\rangle$

²We don't want to multiply the values, just list them in order!

Multi-bit Inner Product

We can work out the inner product using the single bit definitions.

$$\begin{aligned} & \langle ab|cd \rangle \\ &= (\langle a| \otimes \langle b|)(|c \rangle \otimes |d \rangle) \\ &= \langle a|c \rangle \otimes \langle b|d \rangle \end{aligned}$$

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Similar states give 1:

$$\langle 00|00\rangle = 1$$

Different states give 0:

$$\langle 00|11\rangle = 0$$

$$\langle 01|11\rangle = 0$$

$$\langle 10|11\rangle = 0$$

Rules for Operators

However, we can use the tensor product to make larger operators.

$$\begin{aligned}\hat{I} \otimes \hat{N} |11\rangle &= (\hat{I} \otimes \hat{N})(|1\rangle \otimes |1\rangle) \\ &= \hat{I}|1\rangle \otimes \hat{N}|1\rangle \\ &= |1\rangle \otimes |0\rangle \\ &= |10\rangle\end{aligned}$$

Similarly $(\hat{N} \otimes \hat{I}) |11\rangle = |01\rangle$ and $(\hat{N} \otimes \hat{N}) |11\rangle = |00\rangle$

Outer Product

So far we've learned the **tensor product**

$$|0\rangle \otimes |1\rangle = |01\rangle$$

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Can you guess what an **outer product** looks like?

Outer Product

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and the **inner product**

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Can you guess what an **outer product** looks like?

$$|0\rangle \langle 1|$$

Outer Product

The outer product can be used to change a state. We start with $|1\rangle$.

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The outer product can be used to change a state. We start with $|1\rangle$.

$$\begin{aligned} & |0\rangle \langle 1| |1\rangle \\ &= |0\rangle \langle 1|1\rangle \\ &= |0\rangle \times 1 \end{aligned}$$

Outer Product

The outer product can be used to change a state. We start with $|1\rangle$.

$$|0\rangle \langle 1| |1\rangle$$

$$= |0\rangle \langle 1|1\rangle$$

$$= |0\rangle \times 1$$

The state has been changed from $|1\rangle$ to $|0\rangle$!

$$|OUT\rangle \langle IN|$$

Operators from Outer Products

We already saw that $|0\rangle\langle 1|$ changed 1 to 0.

$|1\rangle\langle 0|$ does the opposite, it changes 0 to 1.

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These are the conditions for the NOT gate and \hat{N} operator.

$$\hat{N} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

Operators from Outer Products

Lets see how he \hat{N} operator works.

$$\hat{N} = |0\rangle \langle 1| + |1\rangle \langle 0|$$

$$\hat{N}|0\rangle = (|0\rangle \langle 1| + |1\rangle \langle 0|)|0\rangle$$

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$$\hat{N}|0\rangle = |0\rangle \times 0 + |1\rangle \times 1$$

$$\hat{N}|0\rangle = |1\rangle$$

The AND Operator

Lets build up the \hat{A} operator.

Input	Output	Outer Product
$ 00\rangle$	$ 0\rangle$	$ 0\rangle\langle 00 $

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Input	Output	Outer Product
$ 00\rangle$	$ 0\rangle$	$ 0\rangle \langle 00 $
$ 01\rangle$	$ 0\rangle$	$ 0\rangle \langle 01 $
$ 10\rangle$	$ 0\rangle$	$ 0\rangle \langle 10 $

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$ 01\rangle$	$ 0\rangle$	$ 0\rangle \langle 01 $
$ 10\rangle$	$ 0\rangle$	$ 0\rangle \langle 10 $
$ 11\rangle$	$ 1\rangle$	$ 1\rangle \langle 11 $

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$ 00\rangle$	$ 0\rangle$	$ 0\rangle \langle 00 $
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$$\hat{A} = |0\rangle \langle 00| + |0\rangle \langle 01| + |0\rangle \langle 10| + |1\rangle \langle 11|$$

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$$\hat{A} = |0\rangle \langle 00| + |0\rangle \langle 01| + |0\rangle \langle 10| + |1\rangle \langle 11|$$

Notice that the states on the right have two bits but the states on the left only have one!

The OR Operator

Lets build up the \hat{O} operator.

Input	Output	Outer Product
$ 00\rangle$	$ 0\rangle$	$ 0\rangle \langle 00 $
$ 01\rangle$	$ 1\rangle$	$ 1\rangle \langle 01 $
$ 10\rangle$	$ 1\rangle$	$ 1\rangle \langle 10 $
$ 11\rangle$	$ 1\rangle$	$ 1\rangle \langle 11 $

$$\hat{O} = |0\rangle \langle 00| + |1\rangle \langle 01| + |1\rangle \langle 10| + |1\rangle \langle 11|$$

Quantum States

Postulates of QM

- 1 State Vector
- 2 Time Evolution
- 3 Measurements
- 4 Composite Systems

Postulates of QM

Postulate 1

Any isolated physical system is completely described by a state vector.

A two state quantum system is called a **Qubit**.

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Where $\alpha, \beta \in \mathbb{C}$ (they are complex numbers) ³

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Postulates of QM

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$$1 = |\alpha|^2 + |\beta|^2$$

³where $\chi = a + ib$ is a complex number, $\chi^* = a - ib$.

Inner Product of Different States

Because we have amplitudes for each state $|0\rangle$ and $|1\rangle$ the inner products of two quantum states have values between 0 and 1.

$$|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

$$|\Phi\rangle = \gamma|0\rangle + \delta|1\rangle$$

$$\langle\Psi|\Phi\rangle = \alpha^*\gamma\langle 0|0\rangle + \beta^*\delta\langle 1|1\rangle$$

$$\langle\Psi|\Phi\rangle = \alpha^*\gamma + \beta^*\delta$$

$$0 \leq \langle\Psi|\Phi\rangle \leq 1$$

Postulates of QM

Postulate 1

Any isolated physical system is completely described by a state vector.

This is true not just for qubits, but for any size of quantum system.

$$|\Psi\rangle = c_0 |\psi_0\rangle + c_1 |\psi_1\rangle + \dots + c_n |\psi_n\rangle$$

Where $\{c_0, c_1, \dots, c_n\} \in \mathbb{C}$

$$|c_0|^2 + |c_1|^2 + \dots + |c_n|^2 = 1$$

Postulates of QM

- 1 State Vector
- 2 Time Evolution
- 3 Measurements
- 4 Composite Systems

Postulates of QM

Postulate 2

Evolution of a closed system is by *unitary transformation*.

Unitary here means that the inner product $\langle \Psi | \Psi \rangle$ is unchanged.

$$\hat{U}|\Psi\rangle = |\Phi\rangle$$

$$\langle \Phi | \Phi \rangle = 1$$

Postulates of QM

- 1 State Vector
- 2 Time Evolution
- 3 Measurements
- 4 Composite Systems

Postulates of QM

Postulate 3

Measurements of quantum systems project the system onto one of its states.

For a single qubit state $|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$, a measurement of the qubit will transform it into the state $|0\rangle$ with probability $|\alpha|^2$ or $|1\rangle$ with probability $|\beta|^2$.

Postulates of QM

- 1 State Vector
- 2 Time Evolution
- 3 Measurements
- 4 Composite Systems

Postulates of QM

Postulate 4

The state vector of a composite system is the tensor product of the component systems state vectors.

$$|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

$$|\Phi\rangle = \gamma|0\rangle + \delta|1\rangle$$

$$|\Psi\rangle \otimes |\Phi\rangle = \alpha\gamma|00\rangle + \alpha\delta|01\rangle + \beta\gamma|10\rangle + \beta\delta|11\rangle$$

Quantum Entanglement

Definition

Two objects are entangled when their joint state $|\Psi\rangle$ cannot be expressed as the product of two states $|\psi\rangle \otimes |\psi\rangle$.

$$|\Psi\rangle \neq |\psi\rangle \otimes |\psi\rangle$$

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$$|\Psi\rangle \neq |\psi\rangle \otimes |\psi\rangle$$

Exercise

Show that the state $|\Psi\rangle = \frac{1}{\sqrt{2}} |00\rangle + \frac{1}{\sqrt{2}} |11\rangle$ is entangled by proving there are no values of $\alpha_0, \alpha_1, \beta_0, \beta_1$ such that:

$$|\Psi\rangle = (\alpha_0 |0\rangle + \beta_0 |1\rangle)(\alpha_1 |0\rangle + \beta_1 |1\rangle)$$

Number of Parameters

In the example where we combine two qubits

$$|\Psi\rangle \otimes |\Phi\rangle = \alpha\gamma |00\rangle + \alpha\delta |01\rangle + \beta\gamma |10\rangle + \beta\delta |11\rangle$$

We need 4 terms to describe the state. If we take the tensor product with another qubit:

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Now there are 8 terms!

Generally, the number of terms is 2^N , where N is the number of qubits.

Too many states

The state/operator notation works great for logical circuits. Inputs and Output are states, and gates are operators which transform the states.

It also works really well for representing quantum states.

However, it gets a little hard to work with for larger states.

State Vectors

We can handle complex states more easily if we switch to expressing states using vector notation.

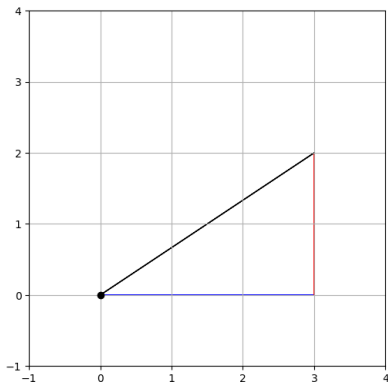
$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Vectors

Vectors are often used to point to a location in space.

For instance, here's a representation of the vector $\begin{pmatrix} 3 \\ 2 \end{pmatrix}$



Bloch Sphere

Vectors can be thought of as arrows with a certain length and direction. Quantum states can be represented by a point on the **Bloch Sphere**.

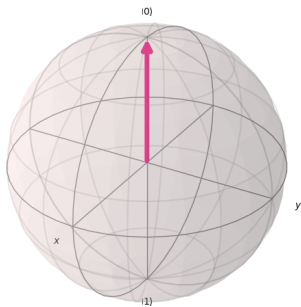


Figure: $|\Psi\rangle = |0\rangle$

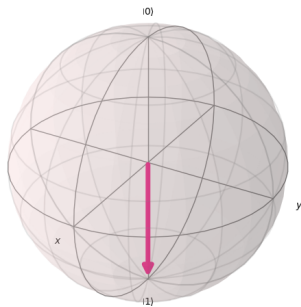


Figure: $|\Psi\rangle = |1\rangle$

We'll explore the the Bloch sphere in the practical parts of the day.

Rules for Vectors

With two vectors $|x\rangle = \begin{pmatrix} a_0 \\ b_0 \end{pmatrix}$ and $|y\rangle = \begin{pmatrix} a_1 \\ b_1 \end{pmatrix}$, we define:

Addition

$$|x\rangle + |y\rangle = \begin{pmatrix} a_0 + a_1 \\ b_0 + b_1 \end{pmatrix}$$

Scalar Multiplication

$$n * |x\rangle = \begin{pmatrix} n * a_0 \\ n * b_0 \end{pmatrix}$$

Inner Product⁴

$$\langle x|y\rangle = a_0^* a_1 + b_0^* b_1$$

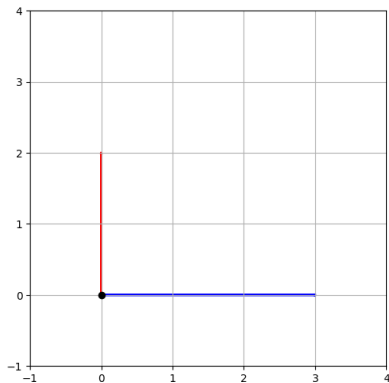
⁴if $z = x + iy$ then $z^* = x - iy$

Vector Addition

The example we saw already illustrated vector addition.

This shows the addition:

$$\begin{pmatrix} 3 \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 2 \end{pmatrix}$$

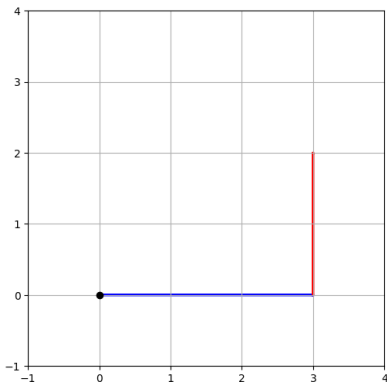


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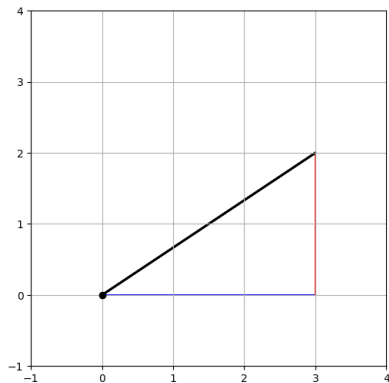


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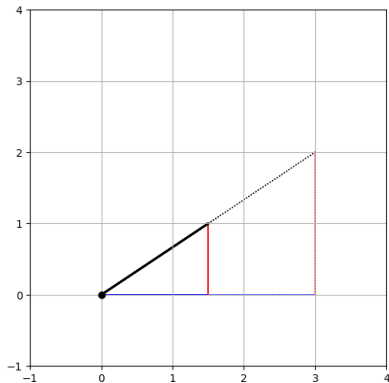
$$\langle x|y\rangle = a_0^* a_1 + b_0^* b_1$$

⁴if $z = x + iy$ then $z^* = x - iy$

Vector Scaling

We can scale this vector too.

$$\frac{1}{2} * \begin{pmatrix} 3 \\ 2 \end{pmatrix} = \begin{pmatrix} 1.5 \\ 1 \end{pmatrix}$$



Rules for Vectors

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$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\langle 0| = (1 \quad 0)$$

$$\langle 1| = (0 \quad 1)$$

Vector Product

The inner product we've been using looks very neat with vectors.

$$\langle 0|0\rangle = (1 \quad 0) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 1 * 1 + 0 * 0 = 1$$

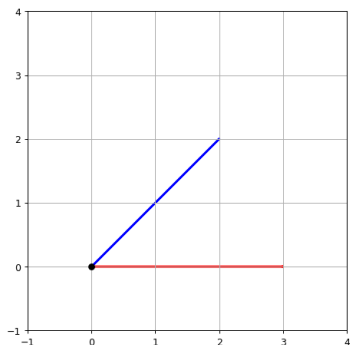
Generally,

$$\langle x|y\rangle = (a \quad b) \begin{pmatrix} c \\ d \end{pmatrix} = ac + bd$$

Inner Product

We can think of the inner product graphically too.

$$\langle x|y \rangle = (3 \ 0) \begin{pmatrix} 2 \\ 2 \end{pmatrix}$$

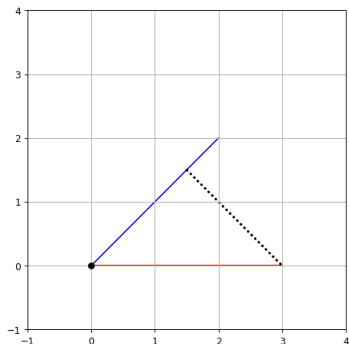


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Then we find the length of the new
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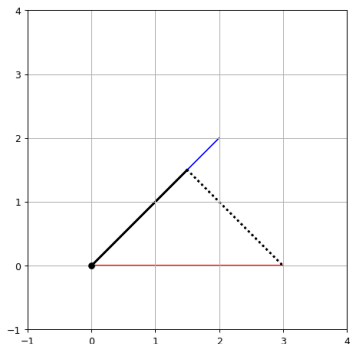
$$\langle x|y \rangle = (3 \ 0) \begin{pmatrix} 2 \\ 2 \end{pmatrix}$$

First we project vector x onto y .

Then we find the length of the new projected vector $|z|$.

Finally, we multiply this by the length of y , $|y|$.

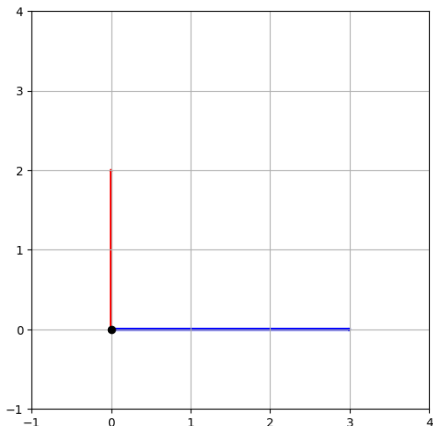
$$\langle x|y \rangle = |z||y|$$



Orthogonal States

By using this projection method, we see that two orthogonal (perpendicular) states have an inner product of 0.

$$\langle x|y \rangle = (3 \ 0) \begin{pmatrix} 0 \\ 2 \end{pmatrix} = 0$$



Computational Basis

Writing our two states $|0\rangle$ and $|1\rangle$ in this way allows us to make any 2 dimensional vector from a combination of them.

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$\begin{pmatrix} a \\ b \end{pmatrix} = a \begin{pmatrix} 1 \\ 0 \end{pmatrix} + b \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

We can this property **spanning** the 2d vectors.

Computational Basis

Additionally, they cannot be written as linear combinations of each other.

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Spanning and linear independence are the two criteria that make a set of vectors a basis.

$\{|0\rangle, |1\rangle\}$ is called the **computational basis** because it relates most clearly to states of bits.

X Basis

Spanning and linear independence seem pretty obvious for the computational basis, but we could make a basis from a different set of vectors.

$$|+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$|-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

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The basis $\{|+\rangle, |-\rangle\}$ is called the X basis, for reasons we'll see soon.

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Exercise

Show that the set $\{|+\rangle, |-\rangle\}$ is a basis for the 2d vectors.

Multi-bit States

To combine single bit states into larger sizes we use a *tensor product*.

$$|x\rangle \otimes |y\rangle = |xy\rangle$$

$$\begin{pmatrix} a_0 \\ b_0 \end{pmatrix} \otimes \begin{pmatrix} a_1 \\ b_1 \end{pmatrix} = \begin{pmatrix} a_0 \begin{pmatrix} a_1 \\ b_1 \end{pmatrix} \\ b_0 \begin{pmatrix} a_1 \\ b_1 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} a_0 a_1 \\ a_0 b_1 \\ b_0 a_1 \\ b_0 b_1 \end{pmatrix}$$

Matrices

A **Matrix** can be thought of as a collection of vectors.

$$M = \begin{pmatrix} a_0 & b_0 & c_0 \\ a_1 & b_1 & c_1 \\ a_2 & b_2 & c_2 \end{pmatrix}$$

Rules for Matrices

Matrices behave quite similarly to vectors.

With two matrices $\hat{X} = \begin{pmatrix} a_0 & b_0 \\ c_0 & d_0 \end{pmatrix}$ and $\hat{Y} = \begin{pmatrix} a_1 & b_1 \\ c_1 & d_1 \end{pmatrix}$

Addition

$$\hat{X} + \hat{Y} = \begin{pmatrix} a_0 + a_1 & b_0 + b_1 \\ c_0 + c_1 & d_0 + d_1 \end{pmatrix}$$

Scalar Multiplication

$$n * \hat{X} = \begin{pmatrix} n * a_0 & n * b_0 \\ n * c_0 & n * d_0 \end{pmatrix}$$

Matrix Multiplication

When a matrix multiplies a vector, it transforms the vector to a new one.

$$\hat{X}|y\rangle = \begin{pmatrix} a_0 & b_0 \\ c_0 & d_0 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \begin{pmatrix} a_0\alpha + b_0\beta \\ c_0\alpha + d_0\beta \end{pmatrix}$$

This just gives us another vector.

Exercise

Find the vector v :

$$v = \begin{pmatrix} 2 & 7 \\ -2 & 1 \end{pmatrix} \begin{pmatrix} -4 \\ 3 \end{pmatrix}$$

Matrix Multiplication

Similarly matrices can be multiplied together

$$\hat{X}\hat{Y} = \begin{pmatrix} a_0 & b_0 \\ c_0 & d_0 \end{pmatrix} \begin{pmatrix} \alpha & \gamma \\ \beta & \delta \end{pmatrix} = \begin{pmatrix} a_0\alpha + b_0\beta & a_0\gamma + b_0\delta \\ c_0\alpha + d_0\beta & c_0\gamma + d_0\delta \end{pmatrix}$$

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Exercise

Find the matrix M :

$$M = \begin{pmatrix} 2 & 7 \\ -2 & 1 \end{pmatrix} \begin{pmatrix} -4 & 0 \\ 3 & 3 \end{pmatrix}$$

Matrix Form of Operators

$$\hat{N}|0\rangle$$

$$\hat{N}|1\rangle$$

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We can actually find out the matrix form of an operator from its outer product form.

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$$\hat{N} = \begin{array}{cc} 0 |0\rangle \langle 0| & 1 |0\rangle \langle 1| \\ 1 |1\rangle \langle 0| & 0 |1\rangle \langle 1| \end{array}$$

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$$\hat{N} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Multi-bit Operators

The **only** time an AND gate returns *ON* is when we have both inputs on.

$$\hat{A}|11\rangle = |1\rangle$$

$$|11\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \otimes \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\hat{A} \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Multi-bit Operators

The \hat{A} operator has the form:

$$\hat{A} = \begin{pmatrix} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

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Multi-bit Operators

and the **only** time an OR gate returns *OFF* is when we have both inputs off.

$$\hat{O}|00\rangle = |0\rangle$$

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$$\hat{O} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

Multi-bit Operators

The $\hat{0}$ operator has the form:

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Logic Summary

We now have three(!) equivalent ways of thinking about computational logic.

- 1 Circuits
- 2 States and Operators
- 3 Vectors and Matrices

These are the same tools we'll use to understand quantum algorithms later on.

Quantum Circuits

Components

- Quantum circuits flow from left to right.

Components

- Quantum circuits flow from left to right.
- We use a wire to represent each qubit.
- Wires cannot be split. (No FANOUT)

Components

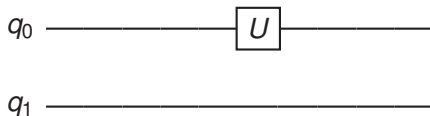
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Pauli Matrices

$$\sigma_X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Pauli Matrices

Once we add in the identity $\mathbb{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$, we have a **basis** for the 2×2 matrices.

Exercise

Show that the set $\{\sigma_x, \sigma_y, \sigma_z, I\}$ is a basis for the 2×2 matrices.

Pauli Matrices

Prove that the Pauli matrices are self-inverse.

$$\sigma_X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\sigma_X \sigma_X = \sigma_Y \sigma_Y = \sigma_Z \sigma_Z = \mathbb{I}$$

Pauli Gates

$$\sigma_X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$



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Hadamard Gate

One of the most common gates is the Hadamard gate, which maximally mixes the $|0\rangle$ and $|1\rangle$ states.

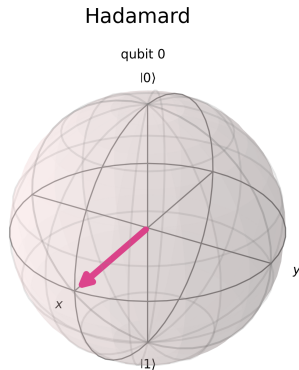
$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

so $\hat{H}|0\rangle = (|0\rangle + |1\rangle)/\sqrt{2}$ and $\hat{H}|1\rangle = (|0\rangle - |1\rangle)/\sqrt{2}$.

Which puts the state half way between $|0\rangle$ and $|1\rangle$.

Visualising Gates

Let's visualise the action of these one qubits gates using the Bloch Sphere.



We can see that the Hadamard gate has moved the state so that it now

Visualising Gates

We know that for a state $|\Psi\rangle = \alpha|0\rangle + \beta|1\rangle$ we expect:

$$\hat{X}|\Psi\rangle = \alpha|1\rangle + \beta|0\rangle$$

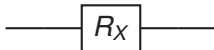
Lets visualise this on the bloch sphere:

Rotation Gates

We can turn these into rotations by an angle θ by mixing them with the identity.

$$R_X = \cos \frac{\theta}{2} \mathbb{I} - i \sin \frac{\theta}{2} \sigma_x$$

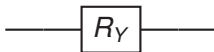
$$R_X = \begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$



Rotation Gates

$$R_Y = \cos \frac{\theta}{2} \mathbb{I} - i \sin \frac{\theta}{2} \sigma_y$$

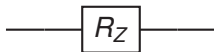
$$R_Y = \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$



Rotation Gates

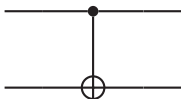
$$R_Z = \cos \frac{\theta}{2} \mathbb{I} - i \sin \frac{\theta}{2} \sigma_z$$

$$R_Z = \begin{pmatrix} \cos \frac{\theta}{2} - i \sin \frac{\theta}{2} & 0 \\ 0 & \cos \frac{\theta}{2} + i \sin \frac{\theta}{2} \end{pmatrix}$$



Two-qubit Gates

Controlled-NOT



$$U_{CN} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

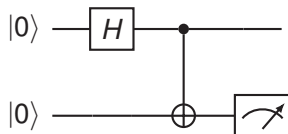
Swap Gate

We can swap the states of two qubits using three CNOT gates.



Measurement

The last basic component is the measurement symbol, which will terminate the line.



IMBQ

Now we're ready to run our first quantum circuits.
We're going to do this using a Jupyter notebook

Metrics

The size of circuits is measured with two different numbers.

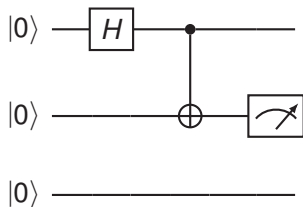
- 1** Depth - The maximum number of gates which have to be applied in sequence.
- 2** Qubit Count - The maximum number of qubits in use at any one time.

Depth

Depth is calculated as the maximum number of gates which must be applied in sequence.

Qubit Count

Qubit count is the maximum number of physical qubits that have to be in use at any one time.



Even though we have three qubits, the qubit count for this circuit is two.

Information & Communication

Information

What sort of things can we do with information?

Information

What sort of things can we do with information?

- Read
- Edit
- Delete
- Move
- Copy
- Encryption / decryption

Readout

When we measure a state we get each possible result with some probability.

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We're interested in knowing these probabilities.

So we have to make lots of measurements.

Which requires preparing the state each time!

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No Cloning Theorem

Theorem

There is no operation which can copy arbitrary quantum states.

We prove this by supposing it is possible.

$$|\Psi\rangle \otimes |s\rangle \xrightarrow{U} |\Psi\rangle \otimes |\Psi\rangle$$

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 (\langle\Psi| \otimes \langle\mathbf{s}|)(|\Phi\rangle \otimes |\mathbf{s}\rangle) &= (\langle\Psi| \otimes \langle\Psi|)(|\Phi\rangle \otimes |\Phi\rangle)
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 \langle\Psi|\Phi\rangle \otimes \langle\mathbf{s}|\mathbf{s}\rangle &= \langle\Psi|\Phi\rangle \otimes \langle\Psi|\Phi\rangle
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 \langle\Psi|\Phi\rangle \otimes \langle s|s\rangle &= \langle\Psi|\Phi\rangle \otimes \langle\Psi|\Phi\rangle \\
 \langle\Psi|\Phi\rangle &= (\langle\Psi|\Phi\rangle)^2
 \end{aligned}$$

So either $|\Psi\rangle$ and $|\Phi\rangle$ are the same state or they're orthogonal!

Information

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Private Key Distribution

Alice and Bob know they will soon have a message to send over the internet, but it's critical it stays private from their friend Eve.⁵

⁵They are planning surprise a party for Eve's cat, who is alive and well. 

Private Key Distribution

Alice and Bob know they will soon have a message to send over the internet, but it's critical it stays private from their friend Eve.⁵

They could meet up in person ahead of time, and agree on a secret **key**, a string of random bits.

$$k = 00111000101100101 \dots$$

⁵They are planning surprise a party for Eve's cat, who is alive and well. 

Private Key Distribution

When the time comes, Alice prepares her message in binary, then adds the key to her own bits modulo 2.

$$\begin{array}{r} 01100010100111011 \\ + \\ 00111000101100101 \\ \hline 01011010001011110 \end{array}$$

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Now her message is encrypted. Bob can decode it by doing the same operation.⁶

Even if Eve intercepted it the message she couldn't know what the original values were without having the key.

⁶in this case addition and subtraction give the same result. Try it out. 

Quantum Key Distribution

This system works great provided that:

- 1 You know you'll need to send the message
- 2 You can meet beforehand
- 3 No one steals either copy of the key

It would be much better to create a new key at the time of each message.

However, if Alice and Bob tried to send each other messages containing the key before the message, Eve could listen in.

Quantum Key Distribution

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- 6 Alice publicly announces b .
- 7 Alice and Bob discard any bits for which b is different from c .

Quantum Key Distribution

Let's see an example.

$a = 0110011$

$b = 1011001$

	a=0	a=1
b=0	$ 0\rangle$	$ 1\rangle$
b=1	$ +\rangle$	$ -\rangle$



$|+\rangle |1\rangle |-\rangle |+\rangle |0\rangle |1\rangle |-\rangle$

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↓

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↓

0011011

Checking for Eavesdropping

The No Cloning Theorem prevents Eve from copying the qubits Alice sends to Bob.

Lets say she'd like to gain information about which state they share without disturbing the state.

$$\begin{aligned} |\psi\rangle |u\rangle &\rightarrow |\psi\rangle |v\rangle \\ |\phi\rangle |u\rangle &\rightarrow |\phi\rangle |v'\rangle \end{aligned}$$

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$$\begin{aligned} \langle \psi | \phi \rangle \langle v | v' \rangle &= \langle \psi | \phi \rangle \langle u | u \rangle \\ \langle v | v' \rangle &= \langle u | u \rangle = 1 \end{aligned}$$

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Because $\langle v | v' \rangle = 1$, they must be the same state. Eve can't learn about Alice and Bob's state without changing it.

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They remove a portion of their key and share publicly to look for discrepancies.

If there are too many differences, they know that someone was intercepting the message.

Hardware

Required properties

The criteria for quantum information processing are:

- Well defined two-level system
- Ability to initialise the state
- Long qubit coherence times⁷
- Universal gate set
- Measurement

⁷Compared to the time it takes to implement a gate.

Desirable properties

These are pretty loose criteria but in reality some designs are better than others.

- Low noise
- Qubit connectivity
- Easy to scale up
- Reliable
- Cheap to build and use

Hardware

- Trapped Ions
- Superconducting Qubits

Trapped Ions

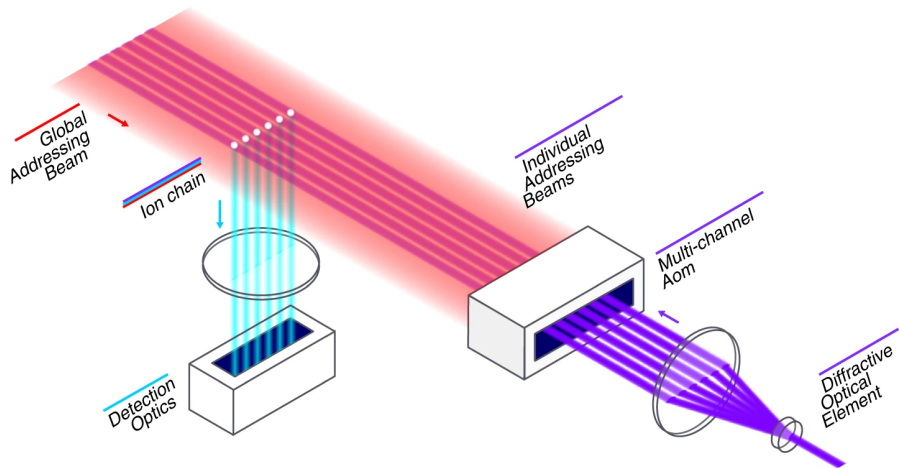


Figure: IONQ

Trapped Ions

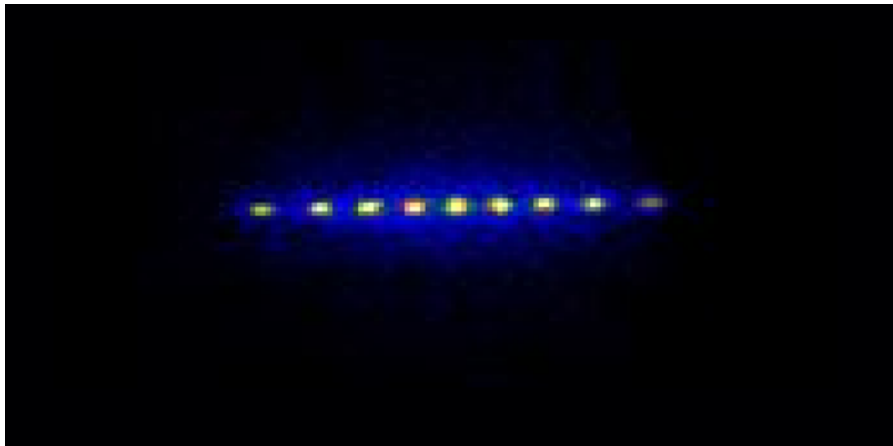
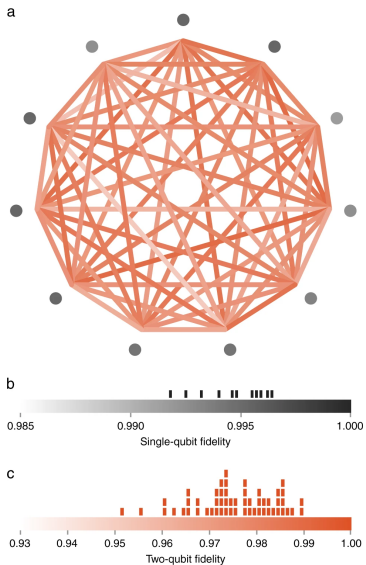


Figure: Trapped Ions - University of Oxford

Trapped Ions



Trapped Ions: Connectivity

Qubits are defined using two energy levels of the ion.

We can manipulate the energy of the ion by using a laser with a resonant frequency.

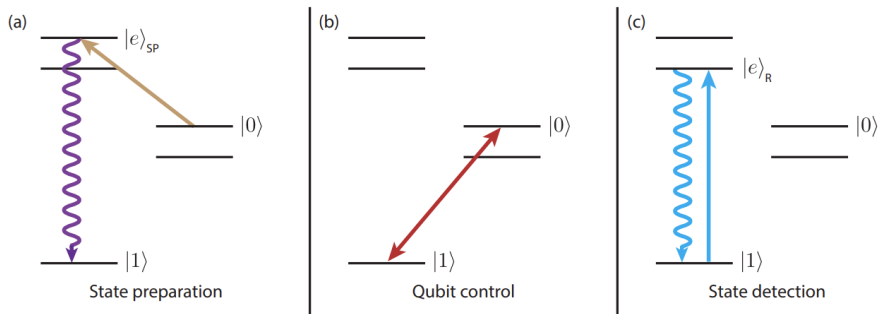


Figure: Bruzewicz et al 2019

Trapped Ions: Gates

The GPI and GPI2 gates are used for single qubit operations.

$$GPI = \begin{pmatrix} 0 & e^{-i\Phi} \\ e^{-i\Phi} & 0 \end{pmatrix}$$

$$GPI2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -ie^{-i\Phi} \\ -ie^{-i\Phi} & 1 \end{pmatrix}$$

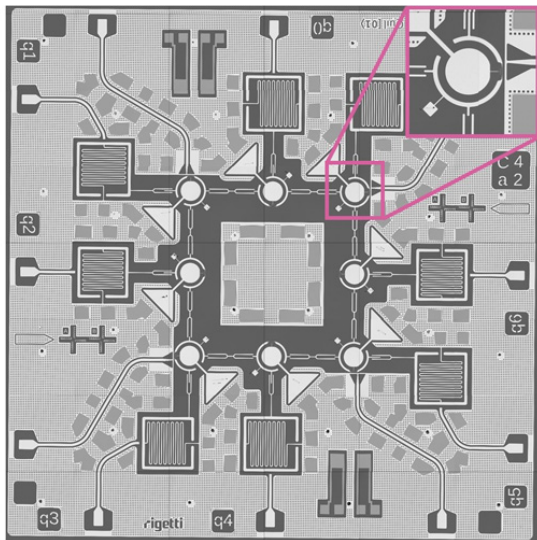
The *Mølmer-Sørensen* gate is used for entangling qubits.

$$\frac{-i}{\sqrt{2}} \begin{pmatrix} i & 0 & 0 & e^{-i(\Phi_0+\Phi_1)} \\ 0 & i & e^{-i(\Phi_0-\Phi_1)} & 0 \\ 0 & e^{-i(\Phi_0-\Phi_1)} & i & 0 \\ e^{-i(\Phi_0+\Phi_1)} & 0 & 0 & i \end{pmatrix}$$

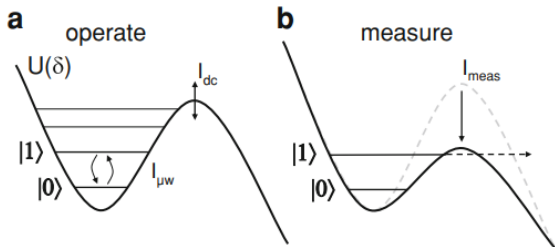
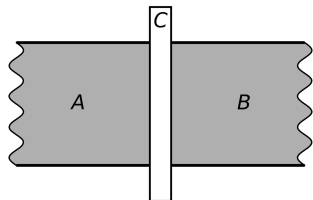
Hardware

- Trapped Ions
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Superconducting Qubits

A

Superconducting Qubits



Superconducting Qubits: Gates

The R_x and R_y gates are used for single qubit operations.

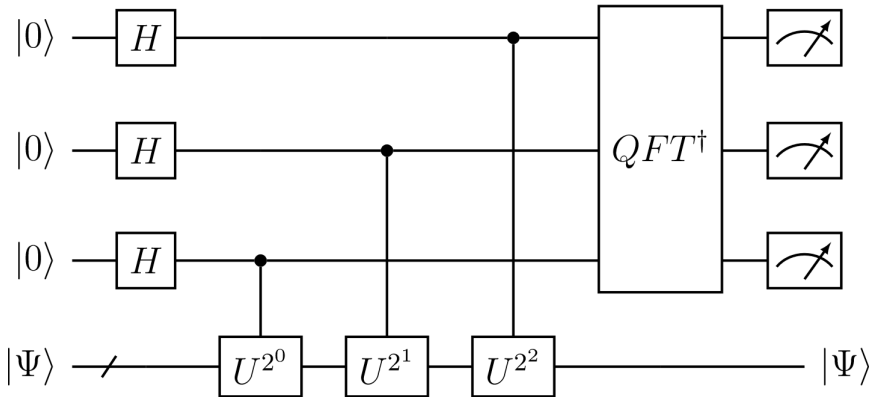
$$R_x = \begin{pmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix} \quad R_y = \begin{pmatrix} \cos \frac{\theta}{2} & -\sin \frac{\theta}{2} \\ \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}$$

The Controlled-Z CZ gate is used for entangling qubits.

$$CZ = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

QPE

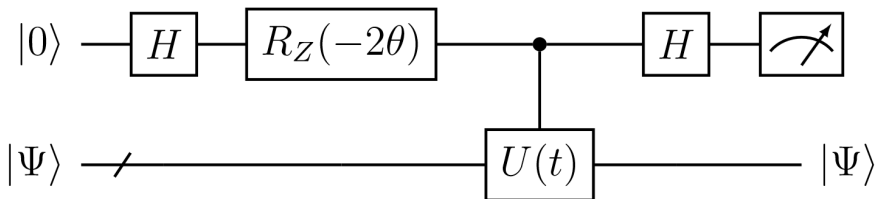
The Quantum Phase Estimator is a way of working out the energy of each state of a system.



Iterative QPE

We can reduce the number of qubits we need by only using one extra at a time.

However, the circuit depth increases.



NISQ

Current quantum computers are very limited in qubit count and depth.

We therefore need to design algorithms for the **Noisy Intermediate Small Quantum Computers**.

Most algorithm research conducted today is focused on NISQ applications.

VQE

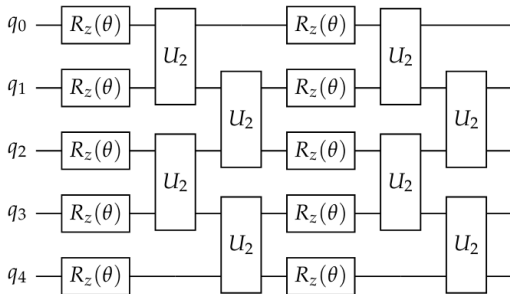
The Variational Quantum Eigensolver is a way to get the lowest energy state $|\psi_0\rangle$ of a system $|\Psi\rangle$.

- 1 Initial state input $|\Phi\rangle$

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- 3 Measure the energy of the state $E = \langle\Psi(\alpha)|\hat{H}|\Psi(\alpha)\rangle$

Postulates of QM

Postulate 1

Any isolated physical system is completely described by a state vector.

This is true not just for qubits, but for any size of quantum system.

$$|\Psi\rangle = c_0 |\psi_0\rangle + c_1 |\psi_1\rangle + \dots + c_n |\psi_n\rangle$$

Where $\{c_0, c_1, \dots, c_n\} \in \mathbb{C}$

$$|c_0|^2 + |c_1|^2 + \dots + |c_n|^2 = 1$$

Variational Principle

The variational principle says that we can find the ground state $|\Psi_0\rangle$ by finding a minimum in the energy of parameterised states.

$$E = \langle \Psi(\alpha) | \hat{H} | \Psi(\alpha) \rangle$$

$$E_0 \leq \langle \Psi(\alpha) | \hat{H} | \Psi(\alpha) \rangle$$

$|\Psi(\alpha)\rangle$ is the ground state $|\Psi_0\rangle$ when we find a minimum in E !

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- 4 Update the parameters $\vec{\alpha}$

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- 2 State evolves according to parameterised circuit $\hat{U}(\vec{\alpha})|\Phi\rangle = |\Psi(\vec{\alpha})\rangle$
- 3 Measure the energy of the state $E = \langle\Psi(\alpha)|\hat{H}|\Psi(\alpha)\rangle$
- 4 Update the parameters $\vec{\alpha}$
- 5 Repeat Steps 1-3 until we find the minimum energy.

VQE Pros & Cons

- Reduces circuit depth.
- Doesn't need any ancilla qubits
- Classical methods for initial state
- Finds real E_0
- Resistant to noise

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- Doesn't need any ancilla qubits
- Classical methods for initial state
- Finds real E_0
- Resistant to noise
- Need to run lots of circuits
- Might not be able to update parameters
- Global or local minimum?

Compilers

Why we need compilers

We can write circuits using lots of languages.

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We want to run algorithms efficiently.

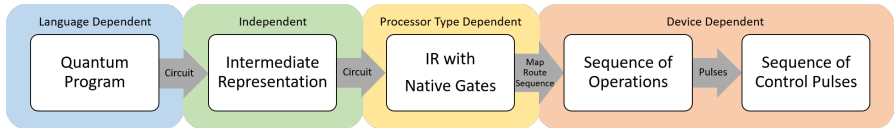
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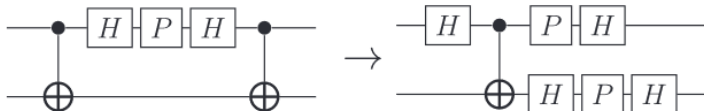
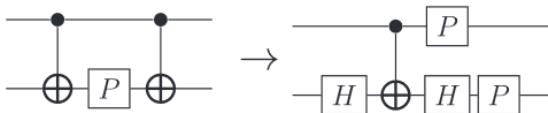
We want to run algorithms efficiently.

Compiler Flow



Redundant Gates

We might inadvertently write a circuit with gates that cancel out.
Compilers remove these using templates.



Intermediate Representation

Programming languages output circuits in a standard format called an **Intermediate Representation**.

Qiskit

Read only

[Open in Quantum Lab](#)

```
1 from qiskit import
   QuantumRegister,
   ClassicalRegister,
   QuantumCircuit
2 from numpy import pi
3
4 qreg_q = QuantumRegister(4,
   'q')
5 creg_c = ClassicalRegister
   (4, 'c')
6 circuit = QuantumCircuit
   (qreg_q, creg_c)
7
8 circuit.h(qreg_q[0])
9 circuit.cx(qreg_q[0], qreg_q
   [1])
10 circuit.measure(qreg_q[1],
   creg_c[1])
```

Intermediate Representation

OpenQASM 2.0

```
1 OPENQASM 2.0;
2 include "qelib1.inc";
3
4 qreg q[2];
5 creg c[2];
6
7 h q[0];
8 cx q[0],q[1];
9 measure q[1] -> c[1];
```

Qiskit

Read only

[Open in Quantum Lab](#)

```
1 from qiskit import
   QuantumRegister,
   ClassicalRegister,
   QuantumCircuit
2 from numpy import pi
3
4 qreg_q = QuantumRegister(4,
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9 circuit.cx(qreg_q[0], qreg_q
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10 circuit.measure(qreg_q[1],
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```

Native Gates

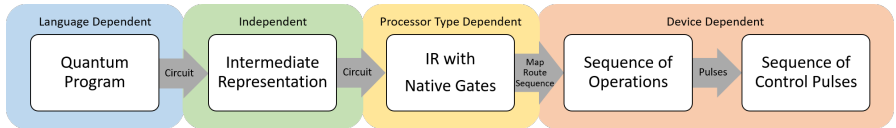
Trapped Ion (IONQ)

- GPI
- GPI2
- Molmer-Sorensen

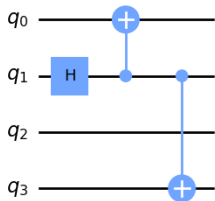
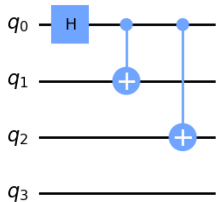
Super Conductor (Rigetti)

- Rx
- Ry
- CZ

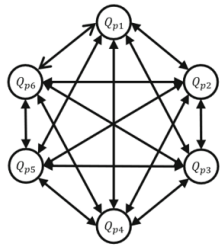
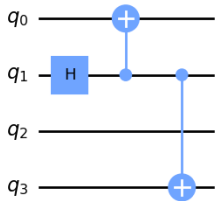
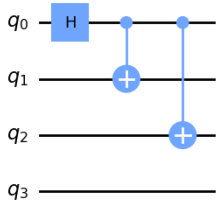
Compiler Flow



Equivalent Circuits



Equivalent Circuits



Connectivity



Figure: Qubit map for IBMQ Manilla

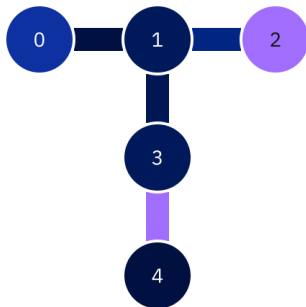


Figure: Qubit map for IBMQ Quito

Quality

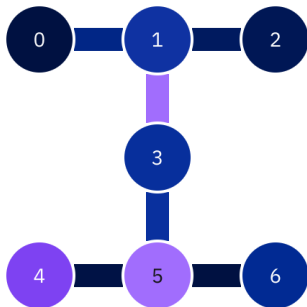


Figure: Qubit map for IBMQ Lagos

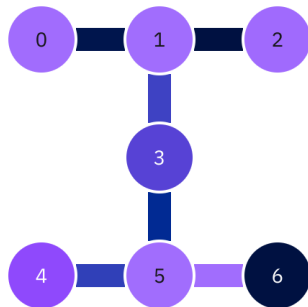
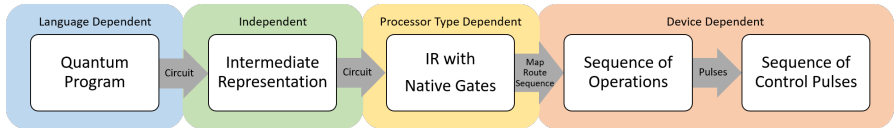


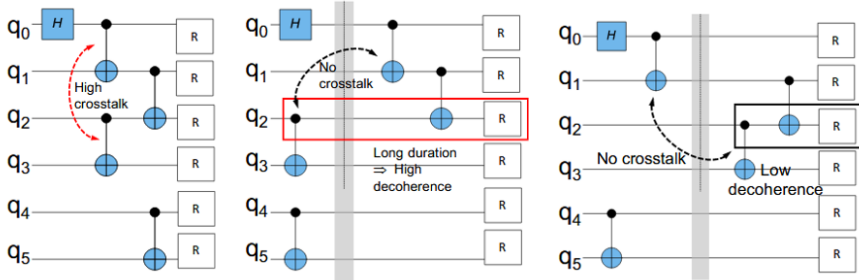
Figure: Qubit map for IBMQ Perth

Compiler Flow



Crosstalk

Crosstalk occurs when we apply two gates at once, causing the signals to interfere.



What we didn't cover

- Adiabatic Quantum Computing
- Non-reversible computing
- Braiding
- Error Correction
- Algorithms

Further Reading

- IMBQ Textbook
- Feynman Lectures
- Nielsen & Chuang